

# BAND STRUCTURE OF MoS<sub>2</sub> USING QUANTUM ESPRESSO

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A PROJECT REPORT

Submitted by

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M MEENAKSHI**

To

The MAHATMA GANDHI UNIVERSITY, KOTTAYAM

In partial fulfilment of the requirements for the award of  
the Degree  
Of  
Bachelor of Science  
In  
Physics



Department of Physics  
Bharata Mata College, Thrikkakara

# DECLARATION

I undersigned hereby declare that the project report "**BAND STRUCTURE OF MoS<sub>2</sub> USING QUANTUM ESPRESSO**", submitted for partial fulfilment of the requirements for the award of degree of Physics of the Mahatma Gandhi University, Kerala is a bonafide work done by me under supervision of Anu Philip , Asst. Professor, Dept. Of Physics. This submission represents my ideas in my own words and where ideas or words of others have been included; I have adequately and accurately cited and referenced the original sources. I also declare that I have adhered to ethics of academic honesty and integrity and have not misrepresented or fabricated any data nor idea or fact or source in my submission. I understand that any violation of the above will be a cause for disciplinary action by the institute and/or the University and can also evoke penal action from the sources which have thus not been properly cited or from whom proper permission has not been obtained. This report has not been previously formed the basis for the award of any degree, diploma or similar title of any other University.

Place: Kochi

DEPARTMENT OF PHYSICS  
BHARATA MATA COLLEGE THRIKKAKARA, KOCHI



## CERTIFICATE

This is to certify that the report entitled BAND STRUCTURE OF MoS<sub>2</sub> USING QUANTUM ESPRESSO submitted by AMRUTHA T AND M MEENAKSHI to the Mahatma Gandhi University in partial fulfilment of the requirements for the award of the Degree of Bachelor of Science in Physics is a bonafide record of the project work carried out by him/her under my/our guidance and supervision. This report in any form has not been submitted to any other University or Institute for any purpose.

Internal Supervisor

External Supervisor

Head of the Department  
Physics

Dr. Anu Philip  
Assistant professor  
Department of Physics

# ACKNOWLEDGEMENT

Before getting into the thickest of things, we would like to thank the personalities who were part of our project in numerous ways, those who gave us outstanding support from the birth of the project. We are extremely thankful to Dr. ANU PHILIP, Assistant Professor of department Of Physics, for providing the necessary infrastructures and resources for the accomplishment of our project and for the esteemed guidance, moral support and invaluable provided by her for the success of the project at BHARATA MATA COLLEGE, KOCHI.

We are highly indebted to DR. JOHNSON K M, the principal of our college, for his support, during the tenure of the project. We are also thankful to all the staff members of the Physics department who have cooperated in making our project a success. We would like to thank all our parents and friends who extended their help, encouragement and moral support either directly or indirectly in our project work.

Thanks for your Valuable Guidance and support.

M. Meenakshi

Amrutha T

# ABSTRACT

In this study, we explore the electronic band structure of molybdenum disulfide (MoS<sub>2</sub>), a prominent two-dimensional material, employing Quantum ESPRESSO, a widely-used package for electronic structure calculations based on density functional theory (DFT). MoS<sub>2</sub> has garnered significant attention due to its unique electronic, optical, and mechanical properties, rendering it promising for various applications in nanoelectronics, optoelectronics, and beyond. By utilizing Quantum ESPRESSO, we calculate the electronic band structure of MoS<sub>2</sub>, which provides crucial insights into its fundamental electronic properties, such as the bandgap, band dispersion, and density of states. Our analysis elucidates the nature of electronic transitions and the presence of band gaps crucial for understanding its behavior in devices. Additionally, we investigate the impact of external perturbations, such as strain or doping, on the band structure, offering valuable guidance for tailoring MoS<sub>2</sub>-based devices for specific applications. Overall, our study contributes to a deeper understanding of the electronic properties of MoS<sub>2</sub>, facilitating its continued exploration and utilization in next-generation electronic and optoelectronic devices.

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# CHAPTER 1

## 1.1 INTRODUCTION

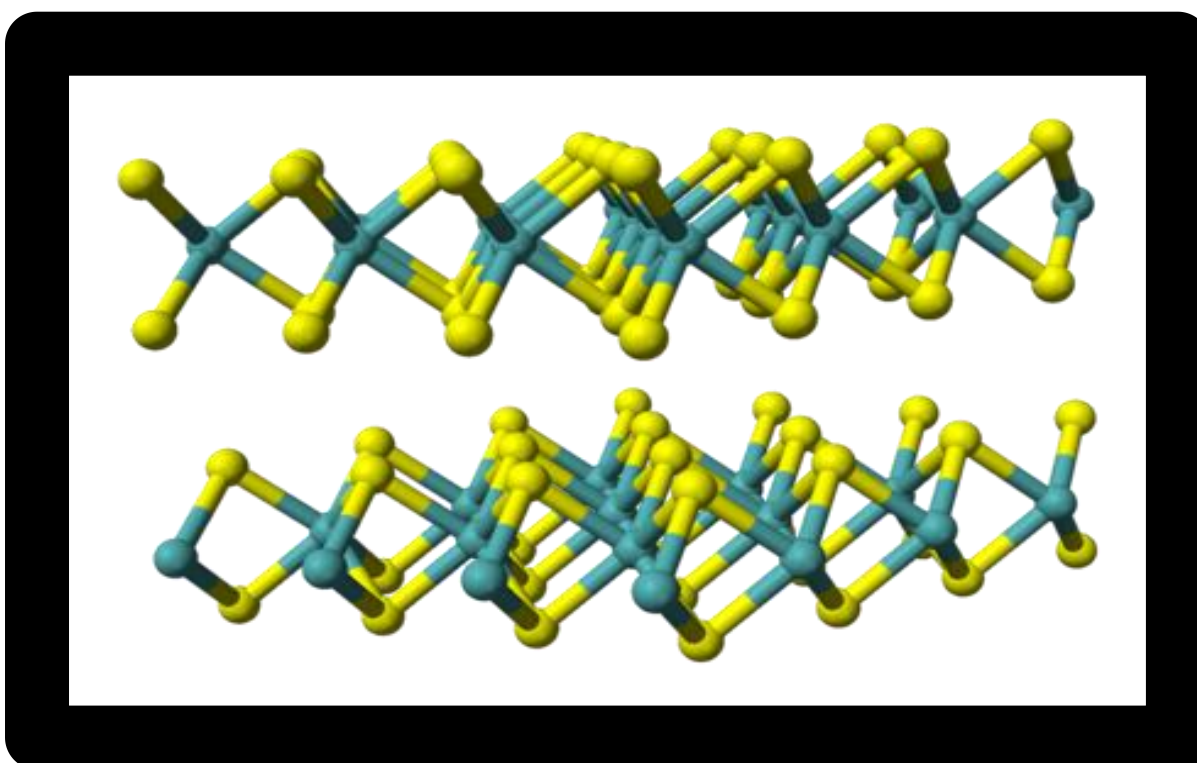
Two-dimensional molybdenum disulfide (2D MoS<sub>2</sub>) has emerged as a remarkable material with transformative properties and versatile applications in various fields of science and technology. Unlike its bulk counterpart, which exhibits an indirect bandgap, monolayer or few-layer MoS<sub>2</sub> possesses a direct bandgap, endowing it with exceptional optical and electronic properties. This unique characteristic makes 2D MoS<sub>2</sub> highly desirable for applications in optoelectronics, catalysis, sensing, and energy storage devices. Its atomically thin structure, coupled with strong quantum confinement effects, facilitates efficient charge carrier transport and manipulation at the nanoscale level. Additionally, the large surface-to-volume ratio of 2D MoS<sub>2</sub> enhances its reactivity and enables precise control over surface chemistry, making it an ideal candidate for catalytic and sensing applications. As research continues to unravel the diverse properties and potential applications of 2D MoS<sub>2</sub>, its significance in advancing nanotechnology and contributing to innovations in various industries becomes increasingly evident.

## 1.2 MoS<sub>2</sub> - BACKGROUND INFORMATION

Molybdenum disulphide (MoS<sub>2</sub>) is an inorganic material composed of 1 molybdenum and 2 sulphide. This chemical has a molar mass of 160.07 g/mol. The substance is categorised as a dichalcogenide of transition metals. The silvery-black solid mineral molybdenite is the primary source of molybdenum. MoS<sub>2</sub> is not very reactive. Oxygen and diluted acids have no effect on it. Graphite and molybdenum disulfide are similar in appearance and feel. Due to its resilience and low friction, it finds extensive application as a dry lubricant. With a bandgap of 1.23 eV, bulk MoS<sub>2</sub> is a diamagnetic, indirect bandgap semiconductor related to silicon.



$\text{MoS}_2$  can be found in nature as either the crystalline mineral molybdenite or the uncommon low-temperature variant of molybdenite called jordisite. Flotation is used to separate the comparatively pure  $\text{MoS}_2$  from molybdenite ore. Carbon is the primary contaminant. In addition, almost all molybdenum compounds can be thermally treated with hydrogen sulphide or elemental sulphur to form  $\text{MoS}_2$ , which can also be created through metathesis events from molybdenum pentachloride.



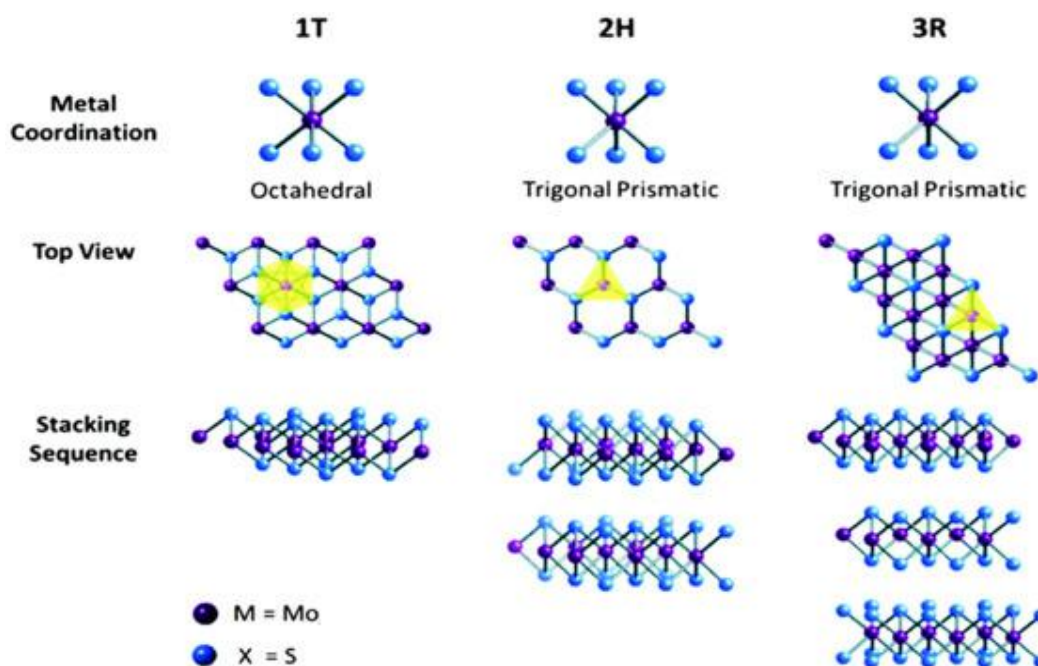


Every form of  $\text{MoS}_2$  is composed of layers, with sulphide ion planes encircling a plane of molybdenum atoms. These three layers combine to form a  $\text{MoS}_2$  monolayer. Weak van der Waals interactions hold stacked monolayers that make up bulk  $\text{MoS}_2$  together.

The names  $2\text{H-MoS}_2$  and  $3\text{R-MoS}_2$ , respectively, refer to two different phases of crystalline  $\text{MoS}_2$ , where the "H" and "R" stand for hexagonal and rhombohedral symmetry. Each molybdenum atom in each of these complexes is covalently bonded to six sulphide ions and is situated at the centre of a trigonal prismatic coordination sphere. Each sulphur atom is coupled to three molybdenum atoms and exhibits pyramidal coordination. The  $3\text{R}$ -phase and the  $2\text{H}$ -phase are both semiconducting. Intercalating alkali metals with  $2\text{H-MoS}_2$  resulted in the discovery of a third, metastable crystalline phase called  $1\text{T-MoS}_2$ . This phase is metallic and possesses trigonal symmetry. The  $1\text{T}$ -phase can be returned to the  $2\text{H}$ -phase by microwave radiation or stabilised by doping with electron donors such as rhenium. By adding S vacancies, the  $2\text{H}/1\text{T}$ -phase transition can be regulated.

$\text{MoS}_2$ , with the chemical formula  $\text{MX}_2$ , belongs to the class of layered 2D transition metal dichalcogenides (TMDCs)[1]. The transition metal element is denoted by M, and S, Se, or Te are represented by X. The stacking of X-M-X layers creates the bulk TMDCs, which can exist in varying structural stages based on the atoms of the transitional metal. Mo (+4) and S (-2) are placed in a S-Mo-S configuration in single-layered  $\text{MoS}_2$  films, with six S atoms surrounding one Mo atom in each layer.  $\text{MoS}_2$  often appears in two structural phases: octahedral ( $1\text{T}$ ) or trigonal prismatic ( $2\text{H}/3\text{R}$ ). The stacking arrangement of layers in  $\text{MoS}_2$  with regard to Mo coordination determines the structural phases. With the ABA layer stacking, the  $2\text{H}$  structure is the predominant stable phase seen in nature. It features two layers with hexagonal symmetry arranged in each unit cell. In this instance, the sulfur atoms from various atomic planes are perpendicular to one another and occupy the same locations. While the S layers are easily able to glide against one other and are connected by van der Waals interactions, the Mo-S connection is primarily covalent.

This structural property helps to cleave crystals along with the surface layers (Z direction) by breaking weak van der Waals forces, and this property is exploited in the synthesis of single-layered  $\text{MoS}_2$



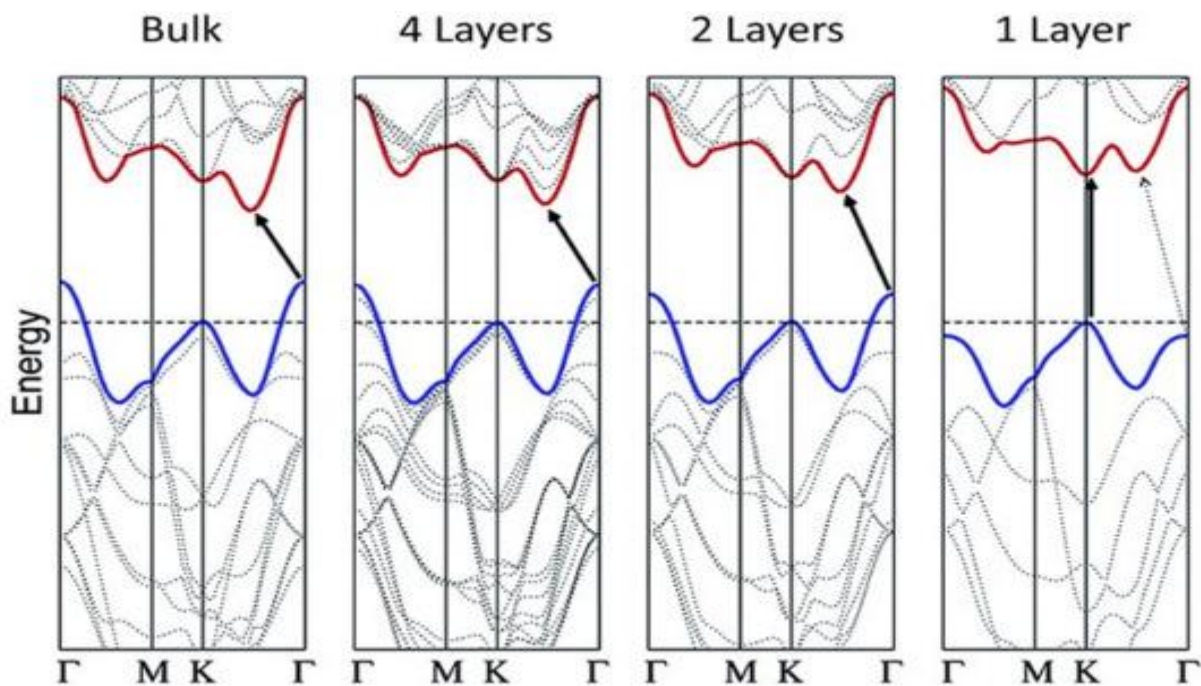
When MoS<sub>2</sub> is synthesised, it often exhibits a 3R structure, which is very unstable and has three layers per unit cell and a rhombohedral symmetry. When heated, this structure's orientation can be changed to a 2H type. The Mo hexagonal arrays are positioned between the sulphur layers in the 2H and 3R MoS<sub>2</sub> structures. One of the sulphur layers in MoS<sub>2</sub> is disoriented to produce the 1T crystal structure, which has an ABC layering sequence. It possesses a trigonal symmetry, or the metallic 1T phase, and a symmetrical Mo–Mo bond. MoS<sub>2</sub> exhibits varying electrical conductivity depending on its phase and crystal structure. Since 1T phases contain more exposed active sites, their presence increases the MoS<sub>2</sub>'s electrical conductivity.

From two-dimensional (2D) sheets of MoS<sub>2</sub> to their three-dimensional (3D) bulk equivalent, the material's band structure varies.

The physical and chemical characteristics of MoS<sub>2</sub> are determined by the number of layers present in the material. While the monolayer MoS<sub>2</sub> has an indirect bandgap of 1.9 eV, the bulk MoS<sub>2</sub> displays a direct bandgap of 1.3 eV. Quantum confinement, interlayer interactions, and long-range Coulomb effects are possible contributors to this phenomena.

In the case of the bulk counterpart, an indirect bandgap is shown between the conduction band (CB) lowest point, K (Brillouin zone point), and the  $\gamma$  (gamma), valence band (VB) maximum point. The CB and VB edges in the illustration are shown by the red and blue lines, respectively. The molybdenum d-orbitals attribute to

the K value, and as this value is independent of the number of layers, it will not alter as the size lowers. Sulphur's 2p orbitals make up the VB. Because it is made up of hybridised orbitals of pz and d orbitals of S and Mo, respectively, the bandgap between K and  $\gamma$  changes as a result. With a decreasing number of layers, the VB maximum (VBM) point  $\gamma$  falls below zero to create a fresh VBM point. The monolayer's lowest energy transition is vertical, or direct, and the energy required for the monolayer's indirect transition is more than that required for the direct band edge transition. There will be an enlarged gap of 1.9 eV in the newly developed band gap between the new  $\gamma$  and the stable K point.



### 1.3 Key Characteristics and Diverse Applications of Two-Dimensional Molybdenum Disulfide (MoS<sub>2</sub>)

- **Layered Structure:** MoS<sub>2</sub> is a two-dimensional substance that resembles graphene in its layered structure. Its structure confers special mechanical, optical, and electrical qualities.
- **Semiconducting Characteristics:** MoS<sub>2</sub> is a semiconductor with a large bandgap, in contrast to graphene, which has a zero bandgap. Because of this characteristic, it may be used in optoelectronic devices, integrated circuits, and transistor applications in electronics.

- **High Carrier Mobility:** MoS<sub>2</sub> demonstrates a critical quality for effective electronic devices: high charge carrier mobility. Because of its great mobility, it can be quickly charged and is hence appropriate for high-speed electronics.
- **Tunable Bandgap:** MoS<sub>2</sub>'s band gap may be adjusted for a variety of optoelectronic and electrical applications by varying its thickness or applying strain.
- **Chemical Stability:** Unlike many other 2D materials that are prone to oxidation, MoS<sub>2</sub> is chemically stable, especially under ambient circumstances. Its usefulness for real-world applications is improved by this steadiness.
- **Huge Surface Area:** MoS<sub>2</sub>'s enormous surface area is a benefit for applications like energy storage, sensing, and catalysis because of its 2-dimensional (2D) structure.
- **Biocompatibility:** Because of MoS<sub>2</sub>'s low toxicity and biocompatibility, it has demonstrated potential in biomedical applications. It can be used for imaging agents, biosensors, and medication delivery systems.
- **MoS<sub>2</sub>'s optical features** are intriguing and include high photoluminescence and nonlinear optical activity. As a result, it may be used in photovoltaic systems, light-emitting diodes (LEDs), and photodetectors.
- **Thermal Conductivity:** MoS<sub>2</sub> has a somewhat high thermal conductivity, however not as high as graphene. This is useful for heat management applications in electronics.
- **Low Cost and Abundance:** The ingredients that make up MoS<sub>2</sub>, sulphur and molybdenum, are inexpensive and abundant when compared to some other 2D materials. This can make large-scale manufacture more financially feasible.

2D MoS<sub>2</sub> exhibits versatile applications due to its absorption spectrum falling within the visible electromagnetic range. Reports indicate its utilization in diverse fields such as energy storage, hydrogen generation, pollutant degradation, and disinfection. Recent studies have extensively documented MoS<sub>2</sub>'s photocatalytic properties across various application domains.

## 1.4 OVERVIEW OF BAND STRUCTURE AND THEIR SIGNIFICANCE IN ELECTRICAL PROPERTIES

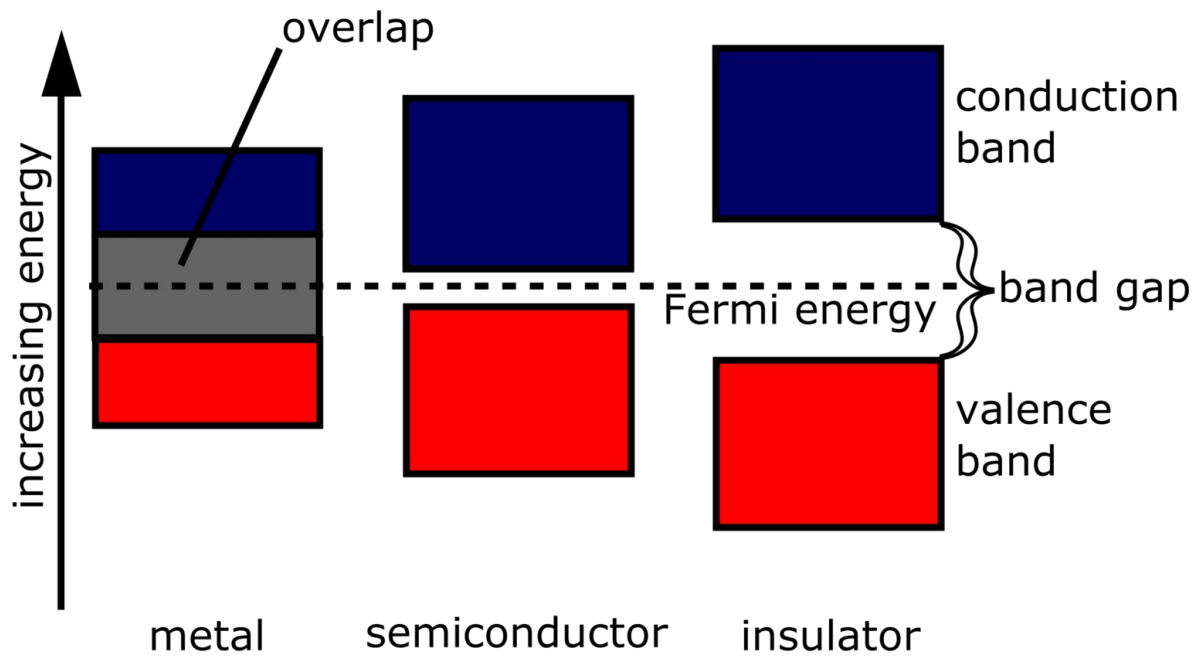
Band structure calculations are pivotal in elucidating the electronic properties of materials, offering crucial insights into their behaviour and functionalities. By examining the band structure, which depicts the distribution of energy levels for electrons within a material, The band structure reveals critical information such as the presence of energy bands, band gaps, and band dispersions, all of which are instrumental in determining the material's conductivity, optical properties, and electronic behaviour. For instance, the presence of a band gap indicates whether a material is an insulator, semiconductor, or metal, thereby defining its electrical conductivity

### 1.4.1 BAND GAP

A band gap, also known as an energy gap or bandgap in solid-state physics, is an energy range in a solid when no electronic states are present. On graphs of the electrical band structure of solids, the term "band gap" refers to the energy difference (usually expressed in electron volts) between the top of the valence band and the bottom of the conduction band in insulators and semiconductors.

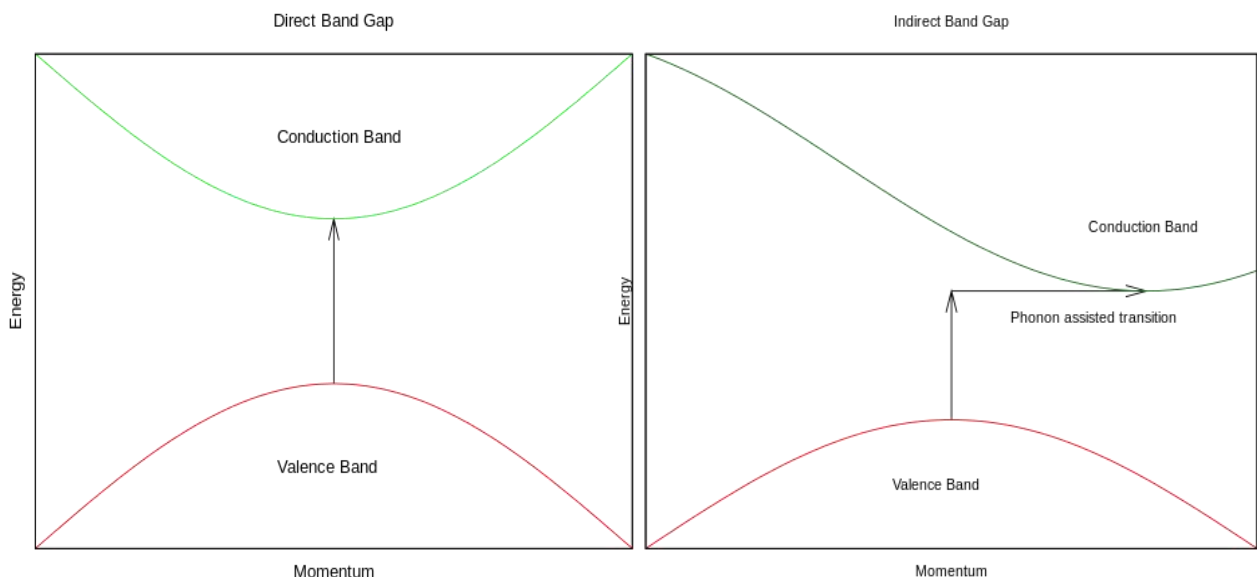
To move an electron from the valence band to the conduction band, energy is needed. In order to carry electric current, the resultant conduction-band electron (as well as the electron hole in the valence band) are free to move inside the crystal lattice. Each type of solid has a unique energy-band structure. The vast diversity of electrical properties seen in different materials can be attributed to this variance in band structure. The spectroscopy and band structure can change depending on the dimension. The valence and conduction bands overlap in conductors. This overlap allows the valence electrons to almost freely enter the conduction band and participate in conduction. A portion of the valence electrons can only pass through the material because of the partial overlap, but this is still sufficient to make conductors conductive.

The gap in semiconductors is small enough to be filled by an excitation of some kind, maybe solar radiation in the case of photovoltaic cells. In essence, the gap is a distance "in-between" an insulator and a conductor.



A semiconductor's band gap can be classified as either an indirect or direct band gap. A certain crystal momentum ( $k$ -vector) in the Brillouin zone characterises the minimal-energy state in the conduction band and the maximal-energy state in the valence band, respectively. The material has an "indirect gap", if the  $k$ -vectors disagree.

When electrons and holes have the same crystal momentum in both the valence band and the conduction band, the band gap is referred to as "direct" because an



electron can emit a photon directly. Since the electron needs to transmit momentum to the crystal lattice through an intermediary state, a photon cannot be released in an "indirect" gap.

Formula for finding direct and indirect band gap of a crystal structure

$$\alpha \propto (h\nu - E_g \pm E_p)^2 \quad \text{in indirect band gap}$$

$$\alpha \propto (h\nu - E_g)^{\frac{1}{2}} \quad \text{in direct band gap}$$

Where alpha is the absorption; h is planck's constant; v is the frequency; E is the energy of the band gap

## 1.5 OBJECTIVES

Various software options are available for analysing and computing the electronic structure and associated parameters of a system. The primary objective of this project is to investigate the bandgap structure and determine the bandgap energy of MoS<sub>2</sub> using Density Functional Theory (DFT) implemented in Quantum Espresso. Initially, the MoS<sub>2</sub> structure is manipulated to meet the necessary conditions for obtaining the band structure. Subsequently, the bandgap energy is calculated from the resulting band structure graphically. This calculated bandgap energy is then compared with theoretically obtained values to assess accuracy. The MoS<sub>2</sub> structure is sourced from the Computational 2D Material Database (C2DB), which hosts over 40 unique crystal structures encompassing various properties of nearly 4000 two-dimensional (2D) materials, including structural, thermodynamic, elastic, electrical, magnetic, and optical properties. Density Functional Theory (DFT) and many-body perturbation theory are employed for computational analysis. Through this investigation, the project aims to contribute to the understanding of MoS<sub>2</sub>'s electronic properties and their potential applications in material science and related fields.

# CHAPTER 2

## 2.1. DENSITY FUNCTION THEORY(DFT)

Density Functional Theory (DFT) revolutionizes material science by providing a robust computational framework for analyzing the electronic properties of materials. At its core, DFT relies on the concept of electron density, representing the spatial distribution of electrons within a material. By solving the Schrödinger equation for electron density, DFT allows researchers to compute critical parameters like total energy, charge density, and wave functions.

DFT finds widespread application in computing band structures, which delineate the distribution of energy levels for electrons in materials. This information offers insights into electrical conductivity and optical properties, facilitating the design of electronic devices and semiconductor materials. Additionally, DFT aids in determining the total energy of materials, crucial for understanding stability, bonding, and phase transitions. By optimizing atomic structures based on total energy minimization, researchers gain valuable insights into material stability and reactivity, with implications for catalysis and materials design.

Furthermore, DFT provides electronic charge density distributions, unveiling details about chemical bonding and inter atomic interactions. By analyzing charge density, researchers discern bonding and non-bonding regions within materials, informing on their chemical and physical properties. Overall, DFT serves as an indispensable tool in materials science, enabling the prediction and understanding of electronic properties in various materials, from semiconductors to metals and insulators. Its versatility and predictive capabilities make DFT instrumental in accelerating materials discovery and advancing applications in fields such as electronics, catalysis, and energy storage



Density Functional Theory (DFT) calculations are instrumental in predicting the electronic properties of materials, offering insights into band structures. DFT computations enable the determination of the electronic band structure of materials. By solving the Kohn-Sham equations within DFT, which describe the behavior of electrons in a material's periodic lattice potential, the energy eigenvalues and corresponding wave functions of electrons are obtained. These results reveal the distribution of energy levels (bands) for electrons in the material, providing critical information about its electronic properties. Band structures elucidate phenomena such as electrical conductivity, optical absorption, and the formation of energy bands and band gaps, which are crucial for understanding and designing electronic and opt-electronic devices.

## 2.2. QUANTUM ESPRESSO

Quantum Espresso is a set of free software available under the GNU General Public Licence that is used for first-principles electronic-structure computations and materials modelling. Density-functional theory (DFT) is its foundation.

An atom, molecule, or condensed phase can be modeled using density-functional theory (DFT), a computational quantum mechanical model used in physics, chemistry, and materials science to investigate the electronic structure (or nuclear structure) of many-body systems (primarily the ground state). With this theory, many-electron systems can be determined by using functionals, i.e. functions of other functions. DFT uses spatially dependent electron densities as functionals. Condensed-matter physics, computational physics, and computational chemistry use DFT as one of their most popular and versatile methods. Solid-state physics has made extensive use of DFT for calculations. Computational materials science utilizes DFT calculations to predict and calculate material behaviour using quantum mechanical considerations, without requiring higher-order parameters like fundamental material properties. A potential acting on the electrons of the system is used in contemporary DFT techniques to evaluate electronic structure.

Hence, Density functional theory is generally used more and more in chemistry and materials science to explain and forecast the behaviour of complex

systems at the atomic scale. DFT computational techniques are specifically used for systems and processing parameters relevant to synthesis. Experimental research in such systems is frequently hindered by non-equilibrium conditions and conflicting results.

Studying the effects of dopants on phase transformation behaviour in oxides, magnetic behaviour in materials containing diluted magnetic semiconductors, and the investigation of magnetic and electronic behaviour in ferroelectrics and dilute magnetic semiconductors are a few examples of modern DFT applications.

BURAI is a GUI system of Quantum ESPRESSO. It is a JavaFX application that allows you to use Quantum ESPRESSO, a powerful and versatile electronic-structure simulation package, in a graphical user interface.

# CHAPTER 3

## CALCULATION SETUP

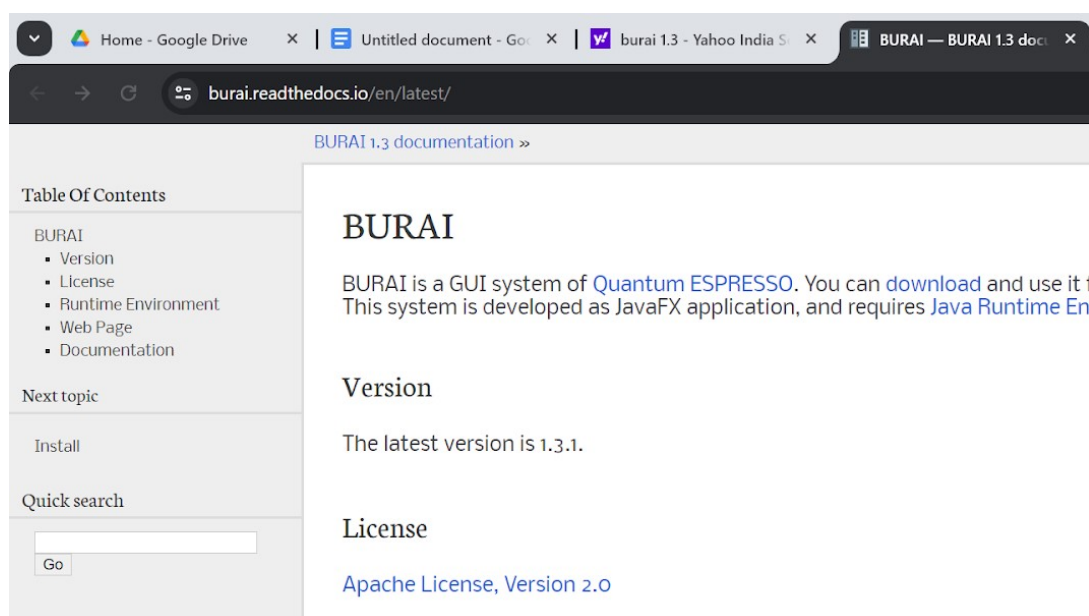
### 3.1. INSTALLING BURAI

STEP 1- Go to '<https://burai.readthedocs.io/en/latest/>' .

STEP 2- install the latest version of Burai i.e. Burai 1.3.1.

STEP 3- If the application depends on the other compressed files in the folder then extract all the files at desired place.

STEP 4- After extracting, run the application. Burai is all set to use.



### 3.2. IMPORTING THE COMPOUND FROM C2DB(Computational 2D Materials Database)

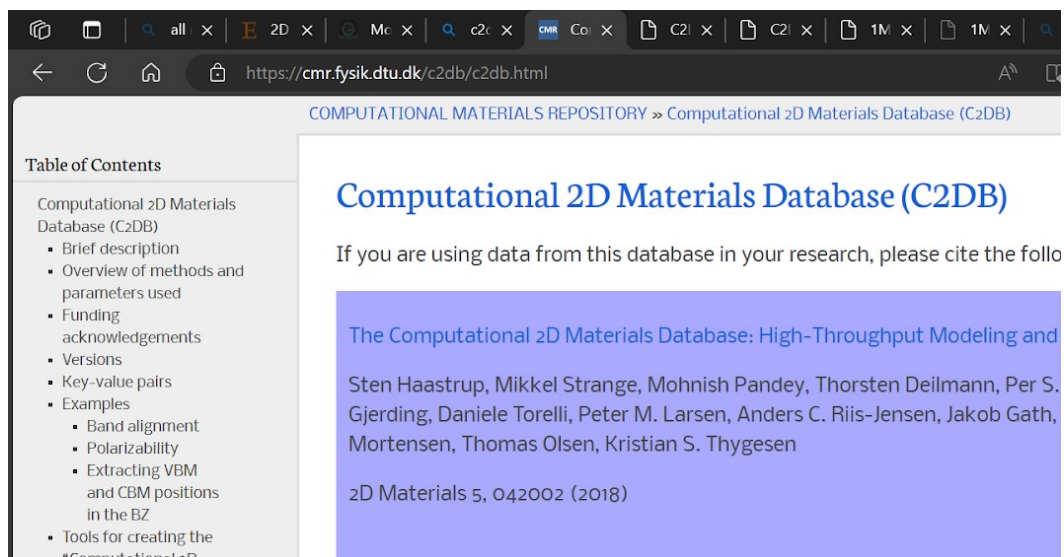
STEP 1- Go to '[Computational 2D Materials Database \(C2DB\) — COMPUTATIONAL MATERIALS REPOSITORY \(dtu.dk\)](#)'.

STEP 2- Go to browse data and search for MOS2.

STEP 3- Select the desired structure from the options.

STEP 4- Download the CIF/XYZ file format of the structure of MoS2.

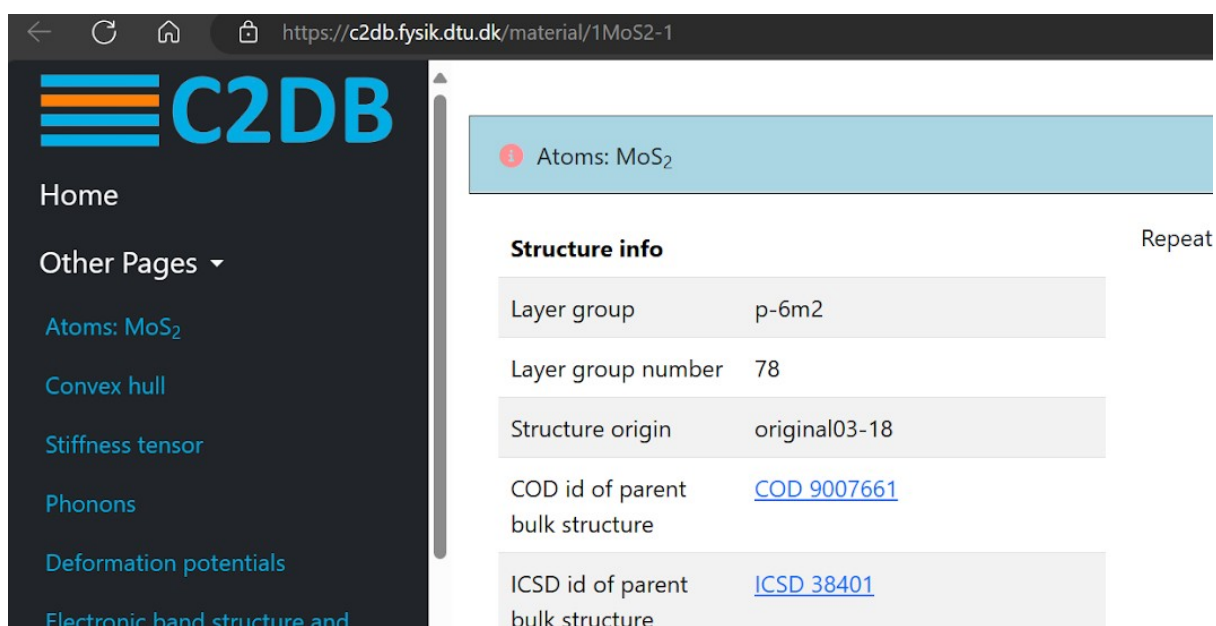
STEP 5- Import the file into the Burai interface.



The screenshot shows the homepage of the Computational 2D Materials Database (C2DB). The browser address bar displays <https://cmr.fysik.dtu.dk/c2db/c2db.html>. The page title is "COMPUTATIONAL MATERIALS REPOSITORY » Computational 2D Materials Database (C2DB)". On the left, there is a "Table of Contents" menu with the following items:

- Computational 2D Materials Database (C2DB)
  - Brief description
  - Overview of methods and parameters used
  - Funding acknowledgements
  - Versions
  - Key-value pairs
  - Examples
    - Band alignment
    - Polarizability
    - Extracting VBM and CBM positions in the BZ
  - Tools for creating the "Computational 2D

The main content area features the title "Computational 2D Materials Database (C2DB)" and a citation request: "If you are using data from this database in your research, please cite the follo". Below this is a blue box containing the citation: "The Computational 2D Materials Database: High-Throughput Modeling and Sten Haastrup, Mikkel Strange, Mohnish Pandey, Thorsten Deilmann, Per S. Gjerding, Daniele Torelli, Peter M. Larsen, Anders C. Riis-Jensen, Jakob Gath, Mortensen, Thomas Olsen, Kristian S. Thygesen 2D Materials 5, 042002 (2018)".



The screenshot shows the material page for MoS<sub>2</sub> in the C2DB. The browser address bar displays <https://c2db.fysik.dtu.dk/material/1MoS2-1>. The page features a dark sidebar with the C2DB logo and a navigation menu:

- Home
- Other Pages ▾
  - Atoms: MoS<sub>2</sub>
  - Convex hull
  - Stiffness tensor
  - Phonons
  - Deformation potentials
  - Electronic band structure and

The main content area shows the material name "Atoms: MoS<sub>2</sub>" with a red warning icon. Below this is a "Structure info" section with a "Repeat" button. The structure information is as follows:

Structure info		Repeat
Layer group	p-6m2	
Layer group number	78	
Structure origin	original03-18	
COD id of parent bulk structure	<a href="#">COD 9007661</a>	
ICSD id of parent bulk structure	<a href="#">ICSD 38401</a>	

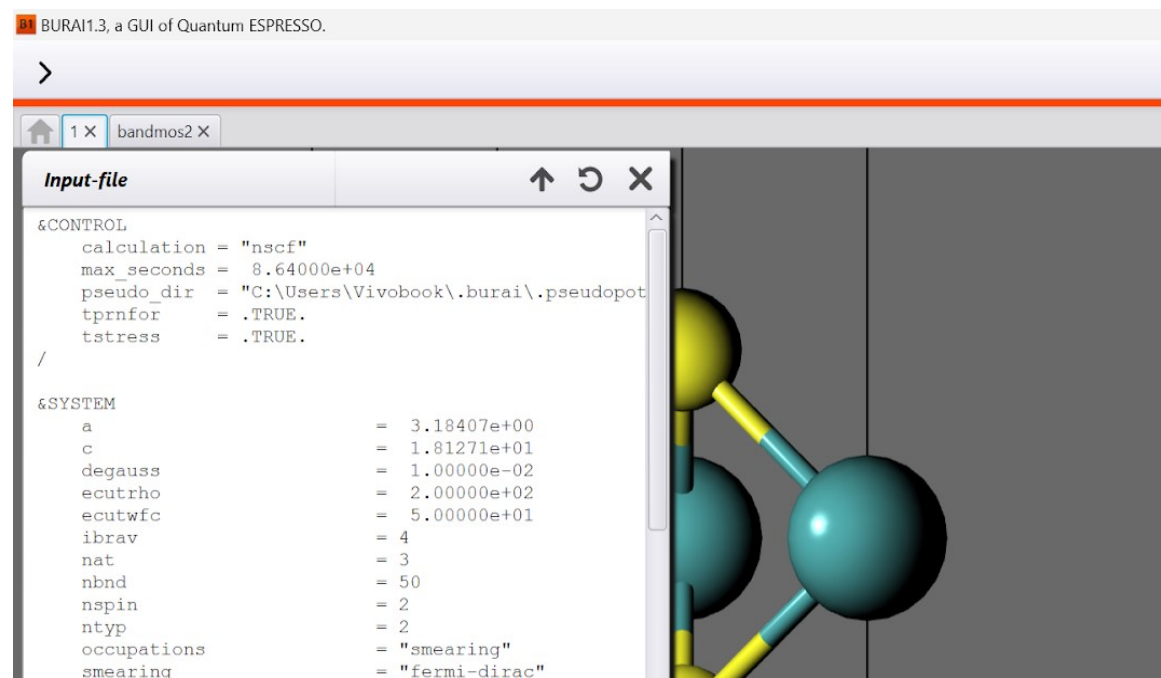
### 3.3.OBTAINING THE BAND STRUCTURE AND CALCULATING THE BANDGAP ENERGY

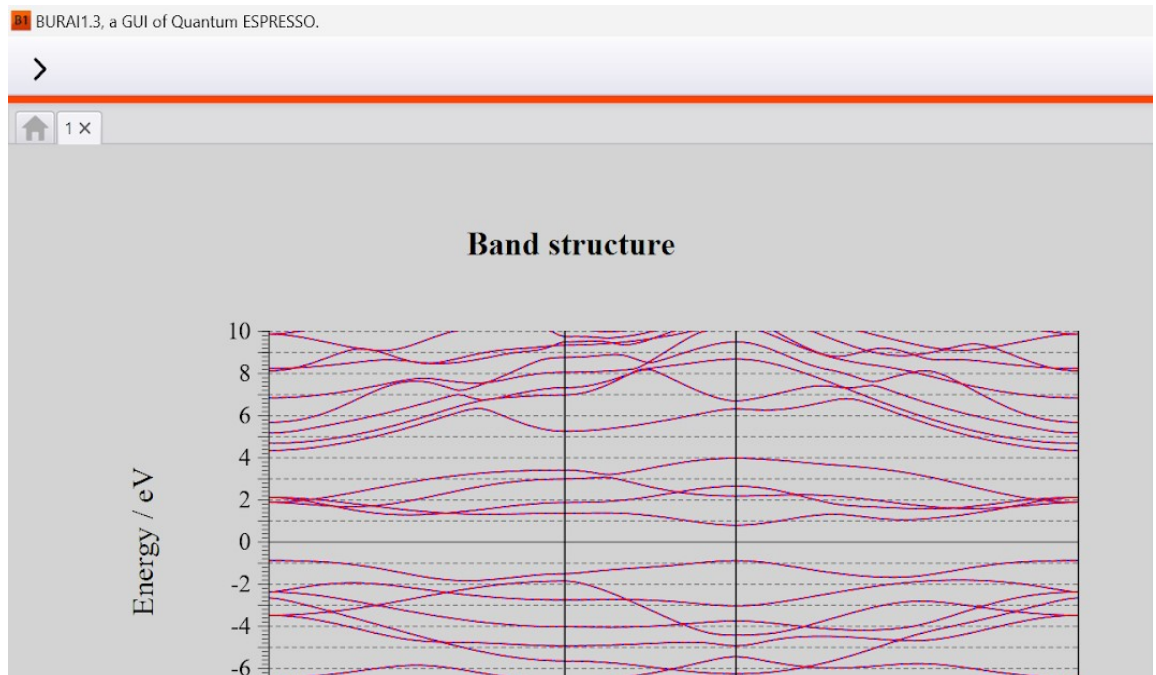
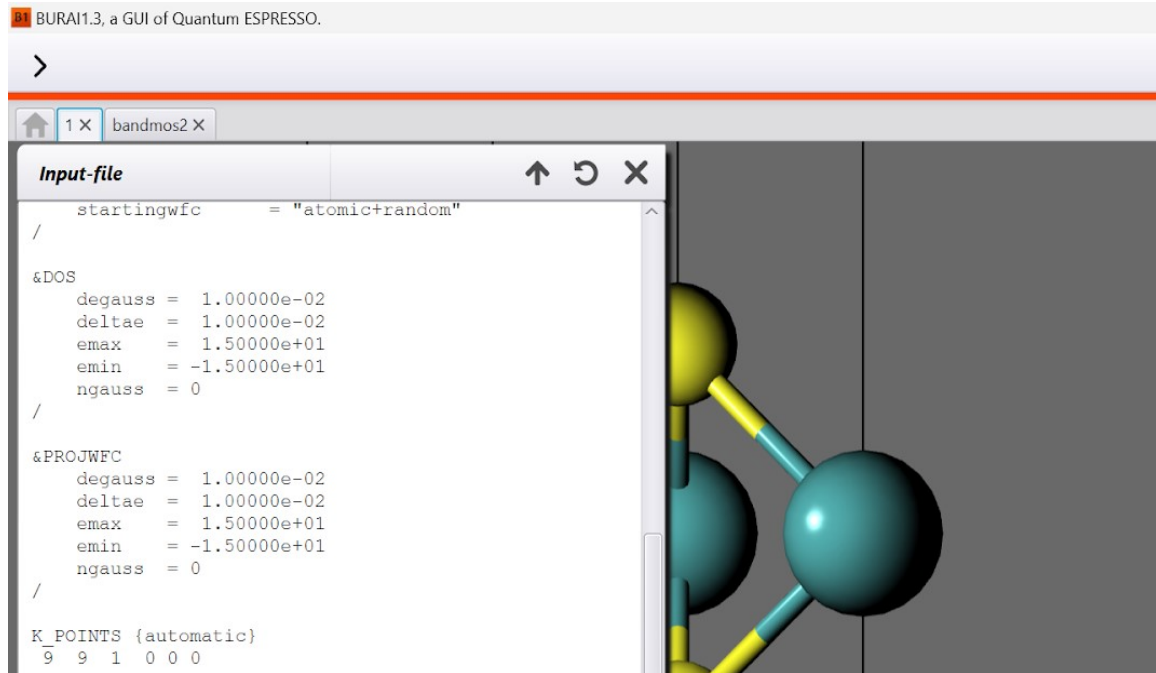
STEP 1- Suitable/required changes are made in the 2D structure of MoS<sub>2</sub> and hence the input file is set up.

**STEP 2-** The structure is saved and run i.e, the calculation of the band takes place.

**STEP 3-** After the calculation is done, the structure output (i.e. the band structure) is obtained under the result tag.

**STEP 4-** The band gap energy is calculated by observing the gap between the lowest conduction band and highest valence band which gives the band gap energy of the structure.





### 3.4. INPUT (DOS PROPERTIES)

&CONTROL

calculation = "bands"

max\_seconds = 8.64000e+04

```

outdir      = "./"
prefix      = "espresso"
pseudo_dir  = "C:\Users\Vivobook\.burai\.pseudopot"
title       = "1(Band)"
tprnfor     = .TRUE.
tstress     = .TRUE.
wf_collect  = .TRUE.
wfcdir     = "./"

```

/

**&SYSTEM**

```

a           = 3.18407e+00
c           = 1.81271e+01
degauss     = 1.00000e-02
ecutrho     = 2.00000e+02
ecutwfc     = 5.00000e+01
ibrav       = 4
nat         = 3
nband       = 50
nspin       = 2
ntyp        = 2
occupations = "smearing"
smearing    = "fermi-dirac"
starting_magnetization(1) = 2.00000e-01
starting_magnetization(2) = 0.00000e+00

```

/

**&ELECTRONS**

```

conv_thr     = 1.00000e-06
electron_maxstep = 200
mixing_beta  = 4.00000e-01
startingpot  = "atomic"
startingwfc  = "atomic+random"

```

/

**&BANDS**

```

filband     = "espresso.band2"
lsym        = .FALSE.
outdir      = "./"
prefix      = "espresso"
spin_component = 2

```

/

**K\_POINTS {tpiba\_b}**

```

4
gG  50
M   50
K   50
gG  50

```



#### ATOMIC\_SPECIES

Mo 95.94000 Mo.pbe-spn-rrkjus\_psl.1.0.0.UPF  
S 32.06600 S.pbe-van\_bm.UPF

#### ATOMIC\_POSITIONS {angstrom}

Mo 0.000000 0.000000 9.063556  
S 1.592033 0.919161 10.627113  
S 1.592033 0.919161 7.500000

### 3.4.1.OPTIMIZATION OF MoS<sub>2</sub>

Title:

1(Optimize)

bravais-lattice index = 4  
lattice parameter (alat) = 6.0170 a.u.  
unit-cell volume = 1074.0408 (a.u.)<sup>3</sup>  
number of atoms/cell = 3  
number of atomic types = 2  
number of electrons = 26.00  
number of Kohn-Sham states= 17  
kinetic-energy cutoff = 50.0000 Ry  
charge density cutoff = 200.0000 Ry  
convergence threshold = 1.0E-06  
mixing beta = 0.4000  
number of iterations used = 8 plain mixing  
Exchange-correlation = SLA PW PBE PBE ( 1 4 3 4 0 0)  
nstep = 201

celldm(1)= 6.017020 celldm(2)= 0.000000 celldm(3)= 5.693060  
celldm(4)= 0.000000 celldm(5)= 0.000000 celldm(6)= 0.000000

crystal axes: (cart. coord. in units of alat)

a(1) = ( 1.000000 0.000000 0.000000 )



a(2) = ( -0.500000 0.866025 0.000000 )

a(3) = ( 0.000000 0.000000 5.693060 )

reciprocal axes: (cart. coord. in units 2 pi/alat)

b(1) = ( 1.000000 0.577350 -0.000000 )

b(2) = ( 0.000000 1.154701 0.000000 )

b(3) = ( 0.000000 -0.000000 0.175652 )

PseudoPot. # 1 for Mo read from file:

C:\Users\Vivobook\.burai\.pseudopot/Mo.pbe-spn-rrkjus\_psl.1.0.0.UPF

MD5 check sum: 5b1023c46552c035e06f08f64a8d82ee

Pseudo is ultrasoft + core correction, Zval = 14.0

Generated using "atomic" code by A. Dal Corso v.5.1.2

Using radial grid of 1229 points, 6 beta functions with:

l(1) = 0

l(2) = 0

l(3) = 1

l(4) = 1

l(5) = 2

l(6) = 2

Q(r) pseudized with 0 coefficients

PseudoPot. # 2 for S read from file:

C:\Users\Vivobook\.burai\.pseudopot/S.pbe-van\_bm.UPF

MD5 check sum: 5981569f150ce26f656387b858196b6a

Pseudo is ultrasoft, Zval = 6.0

Generated by new atomic code, or converted to UPF format

Using radial grid of 811 points, 6 beta functions with:

l(1) = 0

l(2) = 0

l(3) = 1

l(4) = 1

l(5) = 2

l(6) = 2

Q(r) pseudized with 8 coefficients, rinner = 1.100 1.100 1.100  
1.100 1.100

atomic species	valence	mass	pseudopotential
Mo	14.00	95.94000	Mo( 1.00)
S	6.00	32.06600	S ( 1.00)

Starting magnetic structure

atomic species	magnetization
Mo	0.200
S	0.000

12 Sym. Ops. (no inversion) found

Cartesian axes

site n.	atom	positions (alat units)
1	Mo	tau( 1) = ( 0.0000000 0.0000000 2.8465316 )
2	S	tau( 2) = ( 0.4999994 0.2886749 3.3375877 )
3	S	tau( 3) = ( 0.4999994 0.2886749 2.3554759 )

number of k points= 4 Fermi-Dirac smearing, width (Ry)= 0.0100

cart. coord. in units 2pi/alat

k( 1) = (	0.0000000	0.0000000	0.0000000),	wk =	0.0625000
k( 2) = (	0.0000000	0.2886751	0.0000000),	wk =	0.3750000
k( 3) = (	0.0000000	-0.5773503	0.0000000),	wk =	0.1875000
k( 4) = (	0.2500000	0.4330127	0.0000000),	wk =	0.3750000

Dense grid: 51365 G-vectors FFT dimensions: ( 27, 27, 160)

Estimated max dynamical RAM per process > 20.92MB

Estimated total allocated dynamical RAM > 83.70MB

Generating pointlists ...

new r\_m : 0.3126 (alat units) 1.8812 (a.u.) for type 1

new r\_m : 0.3126 (alat units) 1.8812 (a.u.) for type 2

Check: negative/imaginary core charge= -0.000002 0.000000

Initial potential from superposition of free atoms

starting charge 25.99906, renormalised to 26.00000

negative rho (up, down): 9.979E-05 7.607E-05

Starting wfc are 18 randomized atomic wfcs

total cpu time spent up to now is 5.2 secs

### Self-consistent Calculation

iteration # 1 ecut= 50.00 Ry beta=0.40

Davidson diagonalization with overlap

ethr = 1.00E-02, avg # of iterations = 2.5

Threshold (ethr) on eigenvalues was too large:

Diagonalizing with lowered threshold

Davidson diagonalization with overlap

ethr = 6.31E-04, avg # of iterations = 1.2

negative rho (up, down): 1.008E-04 8.578E-05

total cpu time spent up to now is 12.1 secs

total energy = -180.60534820 Ry

Harris-Foulkes estimate = -180.61445141 Ry

estimated scf accuracy < 0.16593173 Ry

total magnetization = 0.02 Bohr mag/cell

absolute magnetization = 0.78 Bohr mag/cell

iteration # 2 ecut= 50.00 Ry beta=0.40

Davidson diagonalization with overlap

ethr = 6.38E-04, avg # of iterations = 2.0

negative rho (up, down): 1.015E-04 9.067E-05

total cpu time spent up to now is 16.1 secs

total energy = -180.62688401 Ry

Harris-Foulkes estimate = -180.63312974 Ry

estimated scf accuracy < 0.07140509 Ry

total magnetization = 0.00 Bohr mag/cell

absolute magnetization = 0.55 Bohr mag/cell

iteration # 3 ecut= 50.00 Ry beta=0.40

Davidson diagonalization with overlap

ethr = 2.75E-04, avg # of iterations = 2.0

negative rho (up, down): 1.023E-04 9.647E-05

total cpu time spent up to now is 19.2 secs

total energy = -180.64606745 Ry  
Harris-Foulkes estimate = -180.64692836 Ry  
estimated scf accuracy < 0.05521308 Ry

total magnetization = 0.00 Bohr mag/cell  
absolute magnetization = 0.48 Bohr mag/cell

iteration # 4 ecut= 50.00 Ry beta=0.40  
Davidson diagonalization with overlap  
ethr = 2.12E-04, avg # of iterations = 1.5

negative rho (up, down): 1.036E-04 1.047E-04

total cpu time spent up to now is 23.3 secs  
total energy = -180.65774984 Ry  
Harris-Foulkes estimate = -180.65248579 Ry  
estimated scf accuracy < 0.00506783 Ry

total magnetization = 0.00 Bohr mag/cell  
absolute magnetization = 0.32 Bohr mag/cell

iteration # 5 ecut= 50.00 Ry beta=0.40  
Davidson diagonalization with overlap  
ethr = 1.95E-05, avg # of iterations = 2.1

negative rho (up, down): 1.028E-04 1.041E-04

total cpu time spent up to now is 27.9 secs

total energy = -180.65900966 Ry  
Harris-Foulkes estimate = -180.65880166 Ry  
estimated scf accuracy < 0.00029637 Ry

total magnetization = 0.00 Bohr mag/cell

absolute magnetization = 0.12 Bohr mag/cell

iteration # 6 ecut= 50.00 Ry beta=0.40

Davidson diagonalization with overlap

ethr = 1.14E-06, avg # of iterations = 2.5

negative rho (up, down): 1.023E-04 1.033E-04

total cpu time spent up to now is 31.5 secs

total energy = -180.65921404 Ry

Harris-Foulkes estimate = -180.65912686 Ry

estimated scf accuracy < 0.00006620 Ry

total magnetization = 0.00 Bohr mag/cell

absolute magnetization = 0.06 Bohr mag/cell

iteration # 7 ecut= 50.00 Ry beta=0.40

Davidson diagonalization with overlap

ethr = 2.55E-07, avg # of iterations = 2.4

negative rho (up, down): 1.024E-04 1.027E-04

total cpu time spent up to now is 37.1 secs

total energy = -180.65926406 Ry

Harris-Foulkes estimate = -180.65924934 Ry

estimated scf accuracy < 0.00000699 Ry

total magnetization = 0.00 Bohr mag/cell

absolute magnetization = 0.03 Bohr mag/cell

iteration # 8 ecut= 50.00 Ry beta=0.40

Davidson diagonalization with overlap

ethr = 2.69E-08, avg # of iterations = 3.0

negative rho (up, down): 1.025E-04 1.026E-04

total cpu time spent up to now is 41.4 secs

total energy = -180.65927312 Ry

Harris-Foulkes estimate = -180.65927273 Ry

estimated scf accuracy < 0.00000124 Ry

total magnetization = 0.00 Bohr mag/cell

absolute magnetization = 0.01 Bohr mag/cell

iteration # 9 ecut= 50.00 Ry beta=0.40

Davidson diagonalization with overlap

ethr = 4.75E-09, avg # of iterations = 3.0

negative rho (up, down): 1.026E-04 1.026E-04

Magnetic moment per site:

atom:	1	charge:	10.7495	magn:	0.0001	constr:	0.0000
-------	---	---------	---------	-------	--------	---------	--------

atom:	2	charge:	4.1520	magn:	-0.0001	constr:	0.0000
-------	---	---------	--------	-------	---------	---------	--------

atom:	3	charge:	4.1520	magn:	-0.0001	constr:	0.0000
-------	---	---------	--------	-------	---------	---------	--------

total cpu time spent up to now is 45.6 secs

End of self-consistent calculation

The Fermi energy is 0.4580 ev

total energy = -180.65927424 Ry  
Harris-Foulkes estimate = -180.65927436 Ry  
estimated scf accuracy < 0.00000027 Ry

The total energy is the sum of the following terms:

one-electron contribution = -963.55685944 Ry  
hartree contribution = 482.37097304 Ry  
xc contribution = -28.50920888 Ry  
ewald contribution = 329.03585576 Ry  
smearing contrib. (-TS) = -0.00003472 Ry

total magnetization = 0.00 Bohr mag/cell  
absolute magnetization = 0.00 Bohr mag/cell

convergence has been achieved in 9 iterations

negative rho (up, down): 1.026E-04 1.026E-04

#### BFGS Geometry Optimization

bfgs converged in 1 scf cycles and 0 bfgs steps  
(criteria: energy < 1.0E-04 Ry, force < 1.0E-03 Ry/Bohr)

End of BFGS Geometry Optimization

Final energy = -180.6592742443 Ry

Begin final coordinates

#### ATOMIC\_POSITIONS (angstrom)

Mo	0.000000000	0.000000000	9.063556000
S	1.592033000	0.919161000	10.627113000
S	1.592033000	0.919161000	7.500000000

End final coordinates



# CHAPTER 4

## 4.1. BAND STRUCTURE

The obtained band structure as an output is as shown,

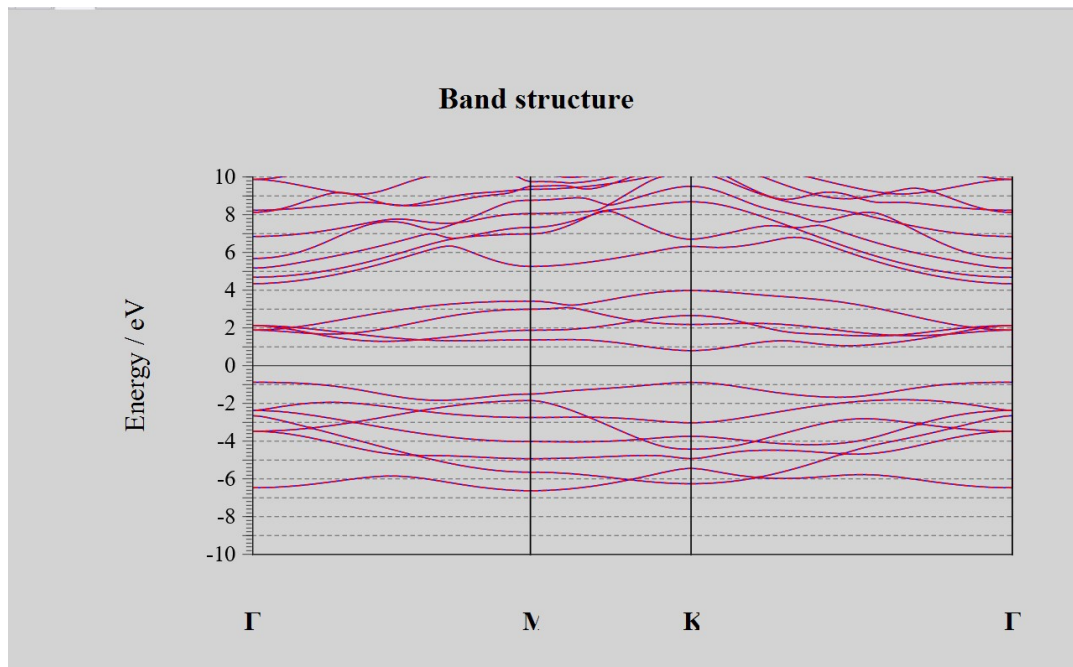


Fig. The band structure of MoS<sub>2</sub> as obtained using quantum espresso.

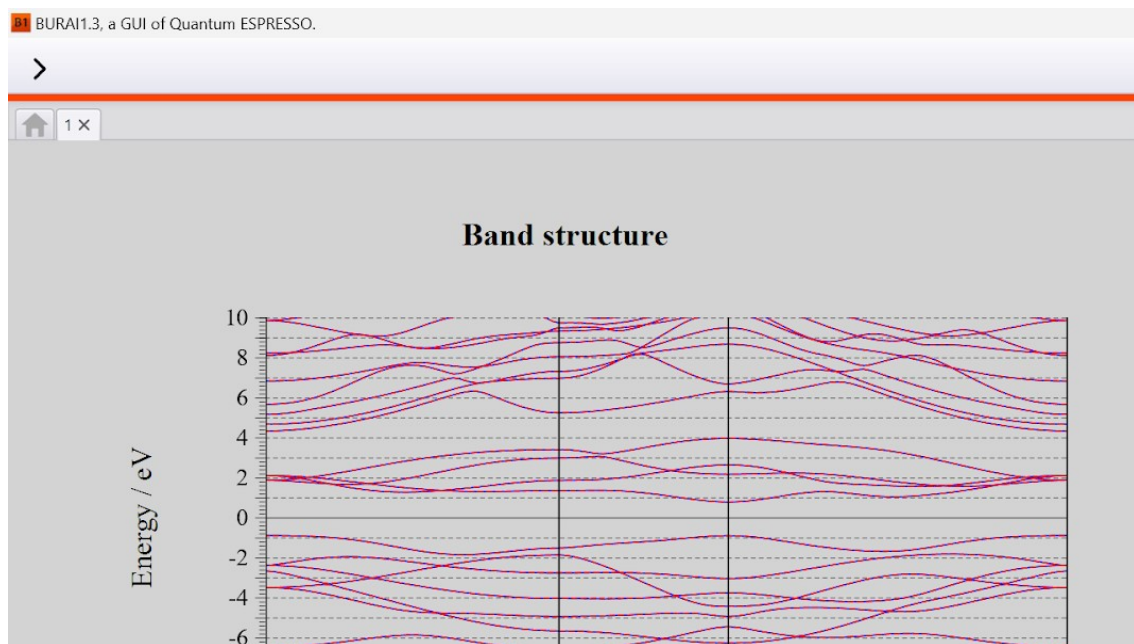


Fig. Output view in the burai interface.

## 4.2. BANDGAP

The band structure of MoS<sub>2</sub> obtained using quantum espresso in the burai interface is a direct band gap from which the band gap energy can be calculated as follows,

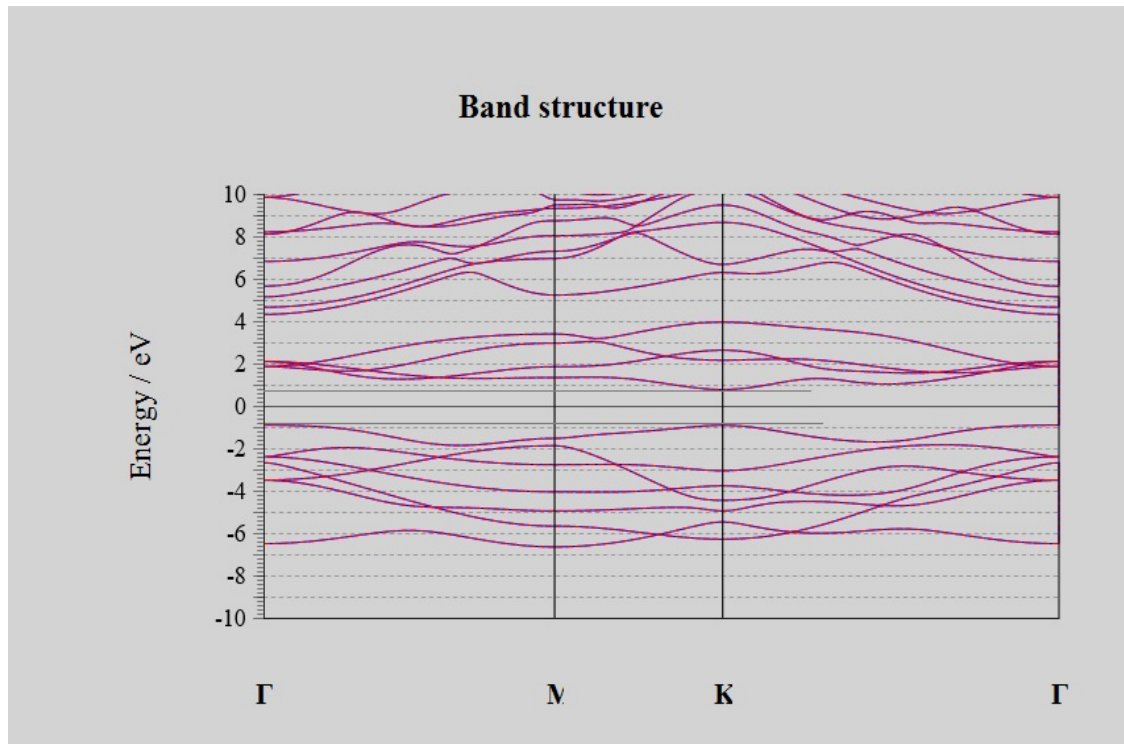


Fig. Band Structure Of MoS<sub>2</sub>.

The band gap is given by the distance between the valence band and conduction band i.e. the energy difference between the top of valence band and bottom of conduction band. It is the energy required to promote an electron from the valence band to the conduction band. From the graph, the value of  $E_v$  is given by

$$E_v = -0.79 \text{ eV}$$

Similarly, the value of  $E_g$  is given by,

$$E_c = 0.8 \text{ eV}$$

The band gap energy is given by energy difference which is given by

$$E_g = E_c - E_v = 0.8 - (-0.79) \\ = 1.59 \text{ eV}$$

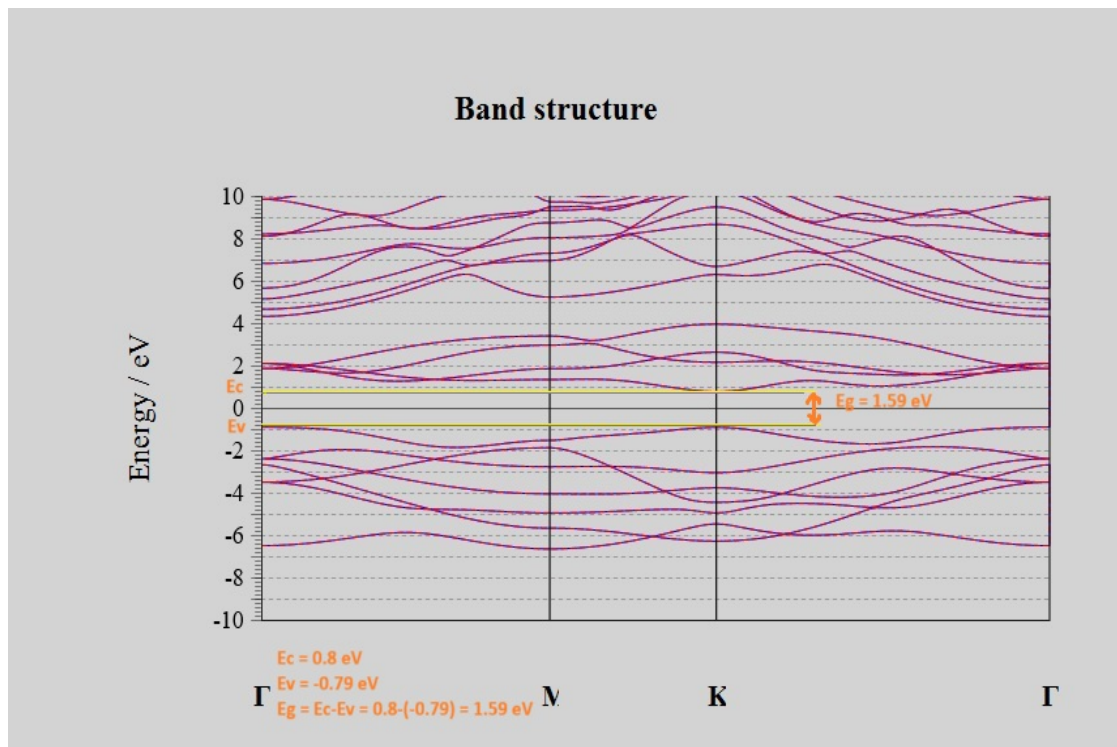


Fig. Calculation of band gap energy from the output band structure.

Therefore, from the graph the band gap energy of the 2D MoS<sub>2</sub> is found to be 1.59 eV.

### 4.3.COMPARISON WITH REPORTED VALUES

The bulk MoS<sub>2</sub> material is reported to have an indirect bandgap of 1.2 eV, whereas two-dimensional (2D) single-layer MoS<sub>2</sub> have a direct bandgap of 1.8 eV(1). From experimental observations and calculations, the bi layer MoS<sub>2</sub> is found to have a bandgap energy of around 1.60 eV(2).

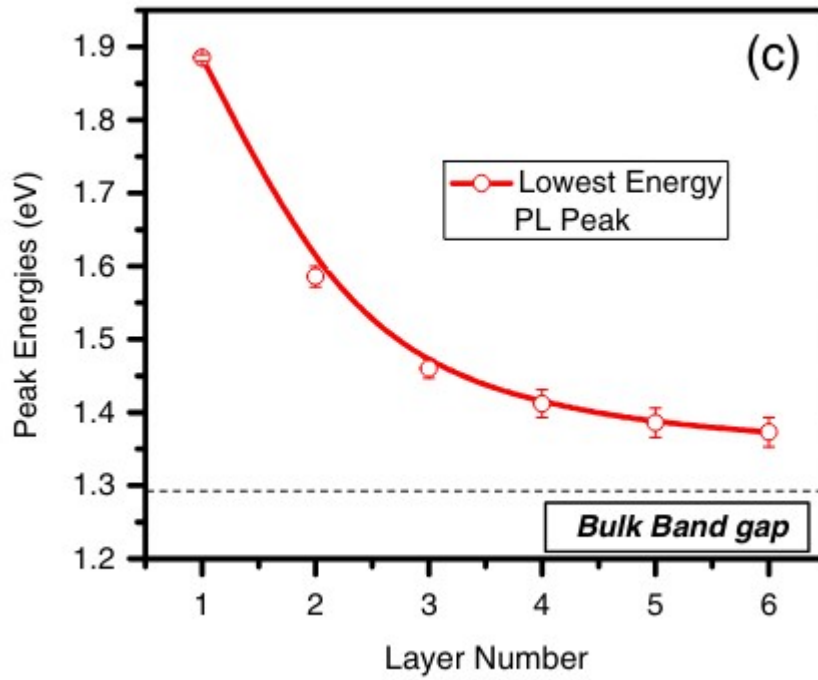


Fig. layer number vs peak energy graph

The experimental band gap energy for the imported bi layer MoS<sub>2</sub> is given to be 1.58 eV.

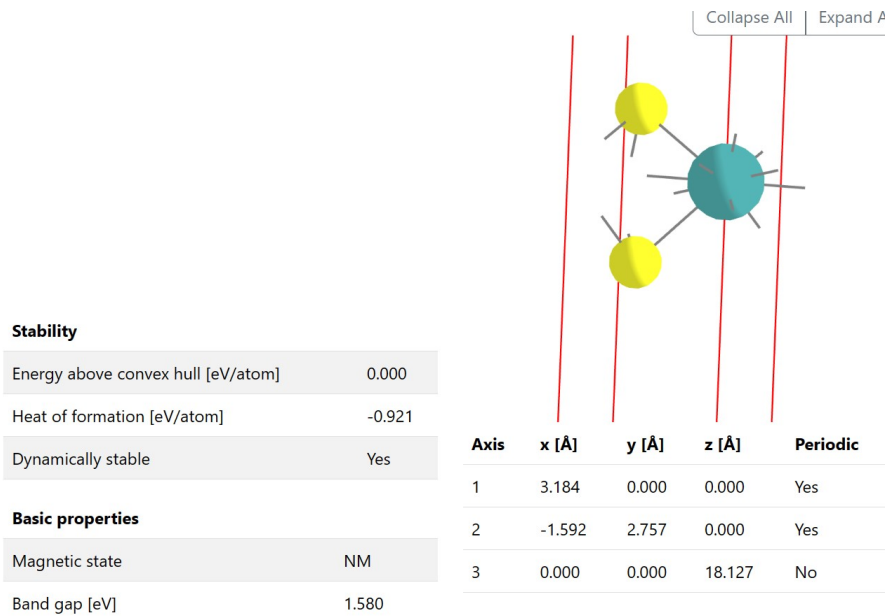


Fig. MoS<sub>2</sub> properties from C2DB.

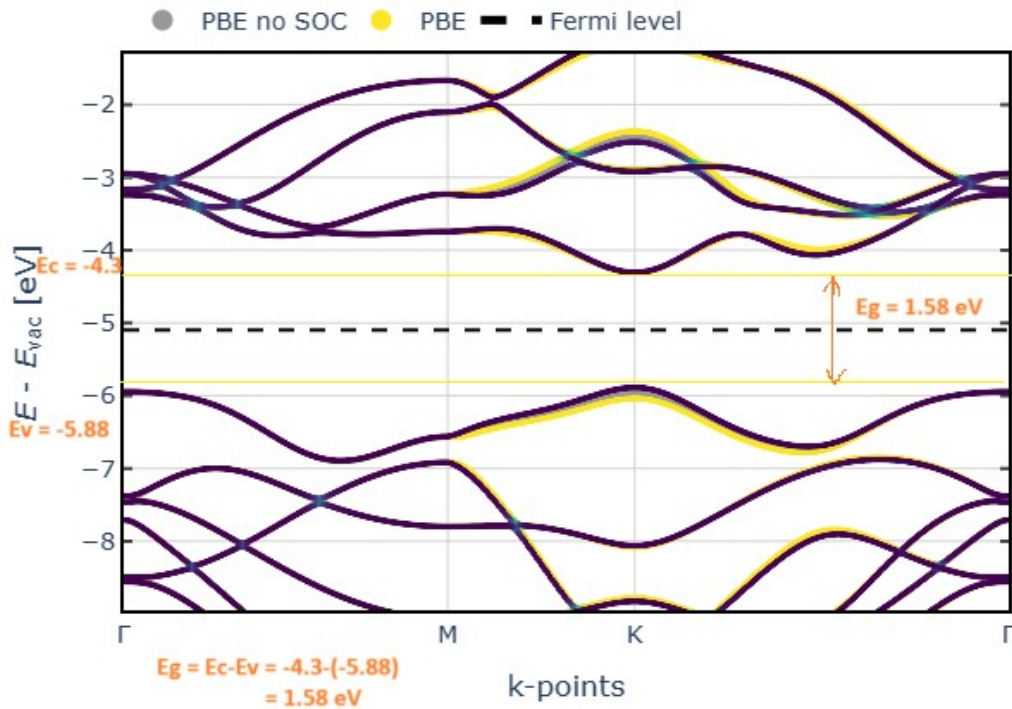


Fig. Direct band structure and calculation of band gap energy of MoS<sub>2</sub> from C2DB.

From the above graph,

$$E_c = -4.3 \text{ eV}$$

$$E_v = -5.88 \text{ eV}$$

Hence the band gap energy is given by,

$$E_g = -4.3 - (-5.88)$$

$$= 1.58 \text{ eV}$$

On Comparison,

Fig. Band gap energy from quantum espresso

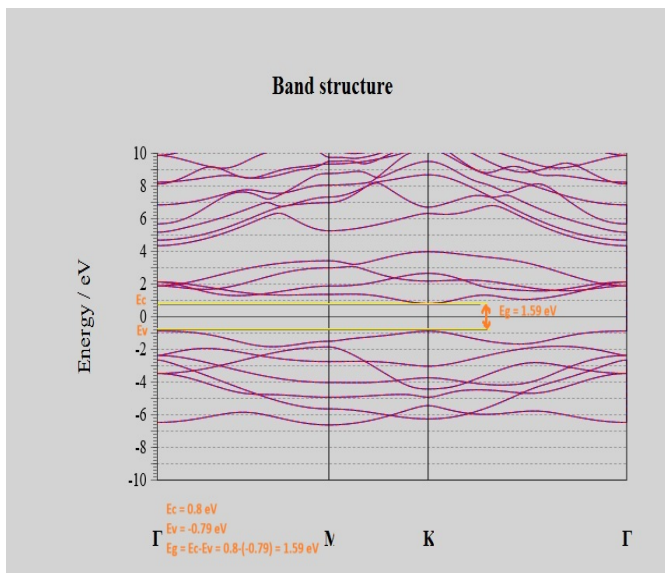
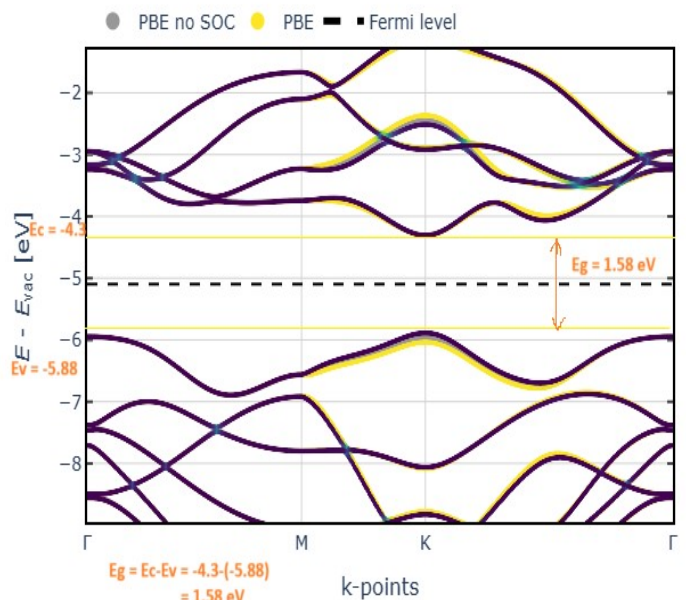


Fig. Band gap energy from C2DB



The band gap energy of MoS<sub>2</sub> theoretically calculated using quantum espresso is found to be 1.59 eV. On considering the above mentioned experimental values and conclusion and hence on further comparison of the value so obtained theoretically and the value obtained experimentally, the band gap energy of MoS<sub>2</sub> can be optimized and interpreted using Quantum espresso.

Tabulating the experimental and calculated values of direct band gap energy of bi layer MoS<sub>2</sub> ,

STRUCTURE	BAND GAP ENERGY (eV)		REFERENCE OF EXPERIMENTAL VALUE
	EXPERIMENTAL	CALCULATED	
BI LAYER	1.58	1.59	<a href="#">1MoS2-1 (dtu.dk)</a>
	1.60		<a href="#">Mak, K. F., Lee, C., Hone, J., Shan, J., &amp; Heinz, T. F. (2010). Atomically ThinMoS2: A New Direct-Gap Semiconductor. Physical Review Letters, 105(13). doi:10.1103/physrevlett.105.136805 .</a>
	1.68		<a href="#">Jiang, J. T., Xiu, S. L., Zheng, M. M., Jia, T. T., Liu, H. Y., Zhang, Y., &amp; Chen, G. (2014). Indirect–direct bandgap transition and gap width tuning in bilayer MoS 2 superlattices. Chemical Physics Letters, 613, 74–79. doi:10.1016/j.cplett.2014.08.060</a>

# CHAPTER 5

## CONCLUSION

Given the experimental data of a 2D material, the structural properties of a compound can be effectively interpreted theoretically. MoS<sub>2</sub> is a vastly explored and scientifically experimented structure in the current world due to its numerous advantages in optical, bio-medical and electronic applications. Band gaps depict the electrical properties of a material signifying whether the material is a metal, insulator or semiconductor. Band gaps refer to the energy difference between the highest occupied electronic state (valence band) and the lowest unoccupied electronic state (conduction band). They play a crucial role in determining the electrical behavior of materials. Materials with a large band gap are insulators, those with a small or no band gap are conductors (metals), and those with a moderate band gap are semiconductors.

MoS<sub>2</sub>, being a semiconductor, exhibits a finite band gap, making it useful in electronic devices where controlled electrical conductivity is required. Additionally, a material might become less reactive and more stable when a band gap is present.

Therefore, it can be concluded that the quantum espresso interface can be effective in generating and obtaining the band structure and thus the band gap Energy of 2D bilayer MoS<sub>2</sub>.

# REFERENCE

[1]

N. Thomas, S. Mathew, K.M. Nair, K. O'Dowd, P. Forouzandeh, A. Goswami, G. McGra naghan, S.C. Pillai. 2D MoS<sub>2</sub>: structure, mechanisms, and photocatalytic applications. <https://doi.org/10.1016/j.mtsust.2021.100073>

[2]

Mak, K. F., Lee, C., Hone, J., Shan, J., & Heinz, T. F. (2010). Atomically Thin MoS<sub>2</sub>: A New Direct-Gap Semiconductor. *Physical Review Letters*, 105(13). doi:10.1103/physrevlett.105.136805.

[3]

Jiang, J. T., Xiu, S. L., Zheng, M. M., Jia, T. T., Liu, H. Y., Zhang, Y., & Chen, G. (2014). Indirect–direct bandgap transition and gap width tuning in bilayer MoS<sub>2</sub> superlattices. *Chemical Physics Letters*, 613, 74–79. doi:10.1016/j.cplett.2014.08.060.

[4]

Z. Y. Zhang, M. S. Si, Y. H. Wang, X. P. Gao, Dongchul Sung, Suklyun Hong, Junjie He Indirect-direct band gap transition through electric tuning in bilayer MoS<sub>2</sub>. *J. Chem. Phys.* 140, 174707 (2014). <https://doi.org/10.1063/1.4873406>

[5]

Computational 2d database (c2db). <https://c2db.fysik.dtu.dk/>